

# Changing Dark Matter

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Based on work done in collaboration with

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# Outline

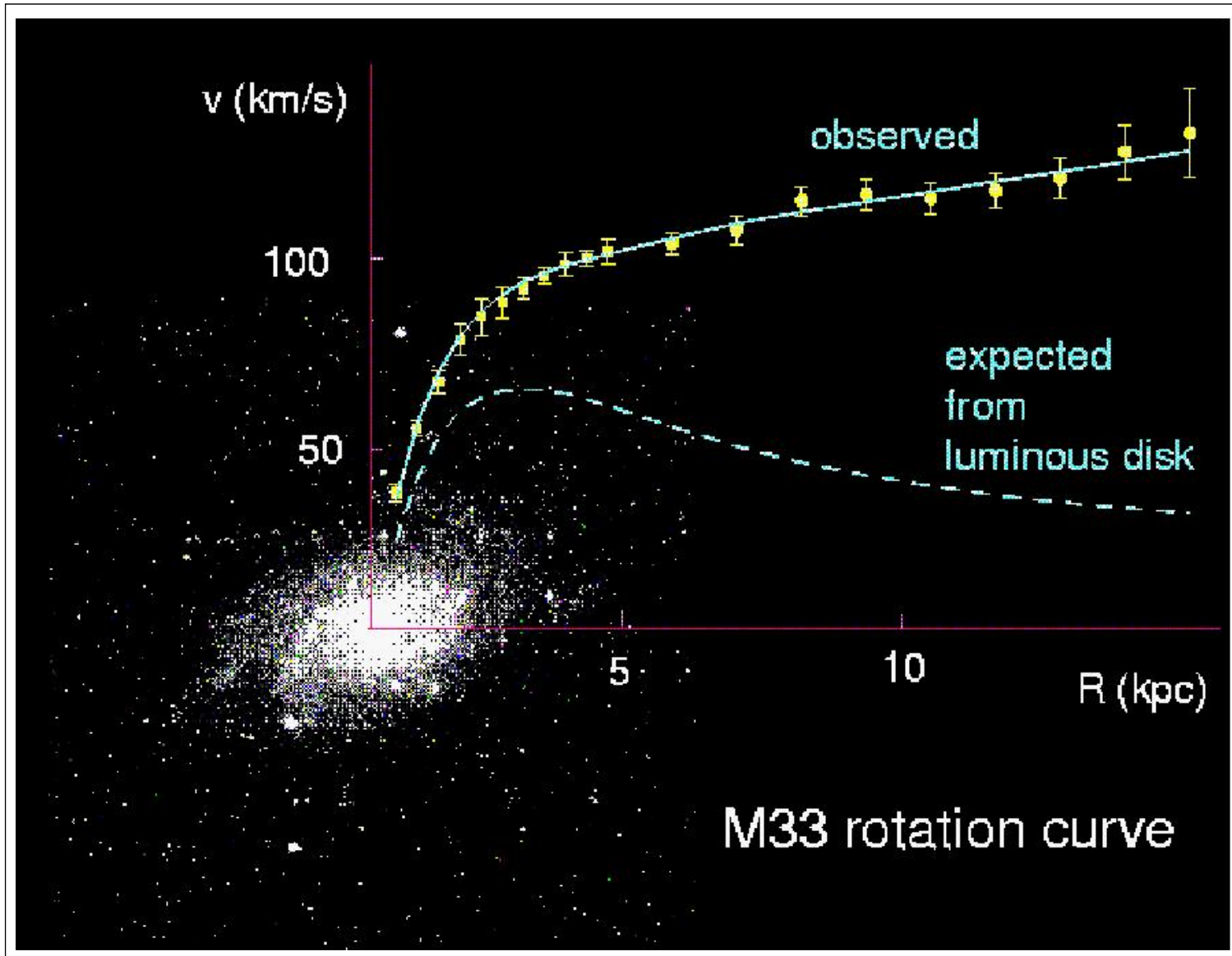
- Evidence and Recent Hints of DM
  - Gravitational Evidence for DM
  - Direct Detection: CDMS, XENON, DAMA, . . .
  - Indirect Detection: PAMELA, ATIC, INTEGRAL, . . .
  - LHC Prospects
- Changing DM
  - A late-time phase transition after freeze-out
  - Effects on the DM
  - Phenomenology of the PT

# Evidence and Recent Hints of Dark Matter

Dark Matter . . .



... not so Dark



# Dark Matter Evidence from Gravity

- WMAP and other CMB probes have precisely determined the average dark matter density in the universe:

$$\Omega_{DM}h^2 = 0.104^{+0.0073}_{-0.0128}.$$

- This value is consistent with estimates of structure formation and gravitational lensing.
- It is also about what we would expect from the thermal freeze-out of a stable WIMP:

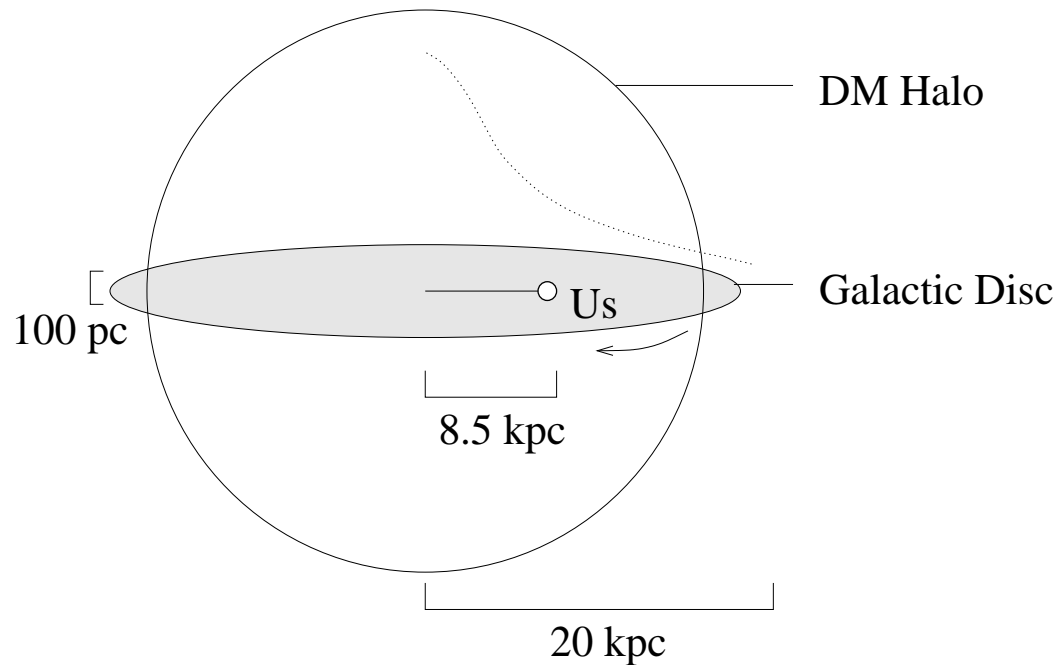
$$\Omega h^2 \simeq \left(\frac{T_0}{M_{\text{Pl}}}\right)^3 \frac{1}{3H_0^2} \left(\frac{m}{T_{fo}}\right) \frac{1}{\langle\sigma v\rangle_{fo}}.$$

## Other Possible Dark Matter Signals

- Dark matter around us can be detected directly by its scattering off nuclei.
- Dark matter in the galaxy can annihilate producing particle fluxes.
- Dark matter captured in the sun can annihilate and generate a neutrino signal.
- If the DM particle is related to a sector that stabilizes the electroweak scale, we may discover it at the LHC.

# DM in our Galaxy

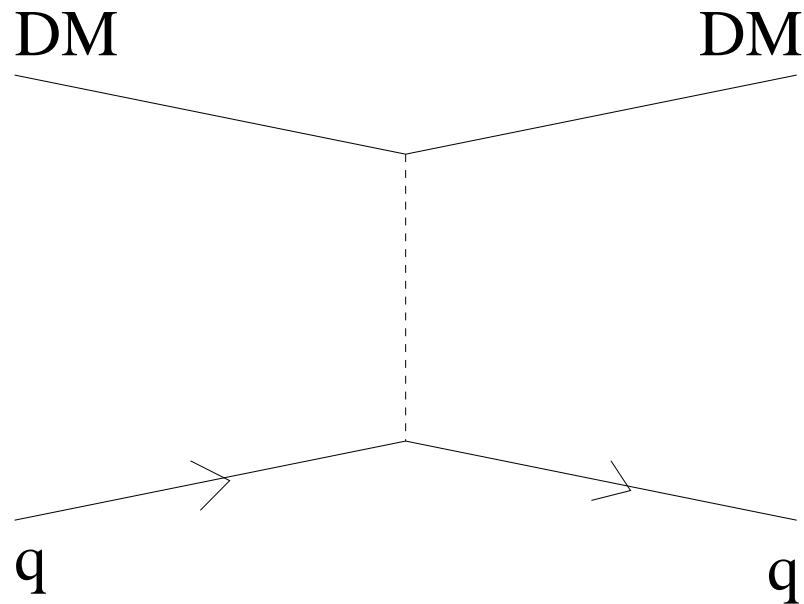
- Flat galactic disc surrounded by a spherical DM halo:



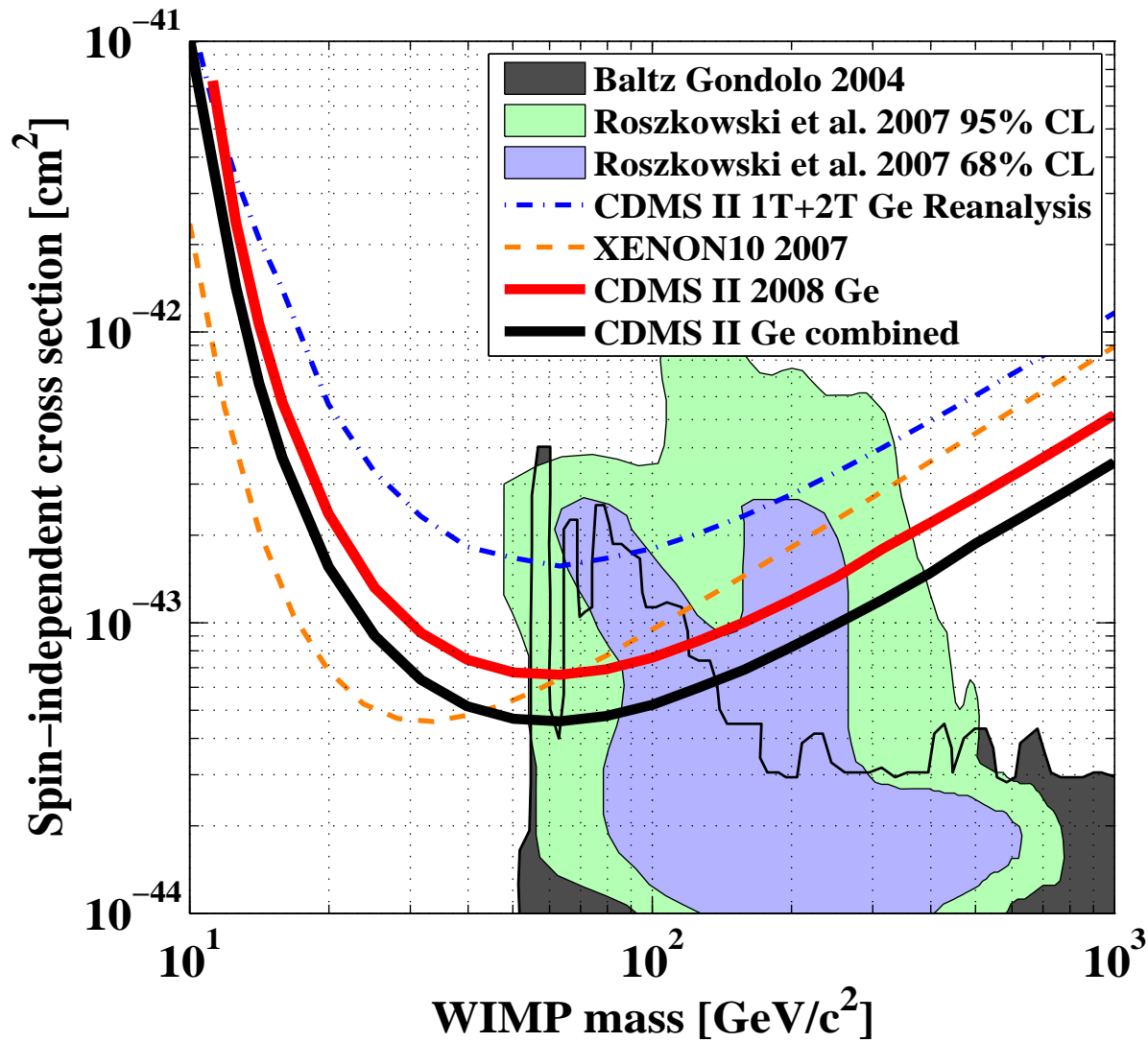
- $v_{us} \simeq (220 \text{ km/s}) + (30 \text{ km/s}) (0.51) \cos(2\pi t/\text{yr})$ .
- $v_{DM} =$  Maxwell distribution with  $\langle v \rangle \simeq 250 \text{ km/s}$ .

# Direct DM Detection

- We encounter a DM “wind” from our motion with the galactic disk.
- This DM flux can scatter off nuclei.  
→ look for nuclear recoils  $\sim 100$  keV



# Experimental Limits (Low Background)



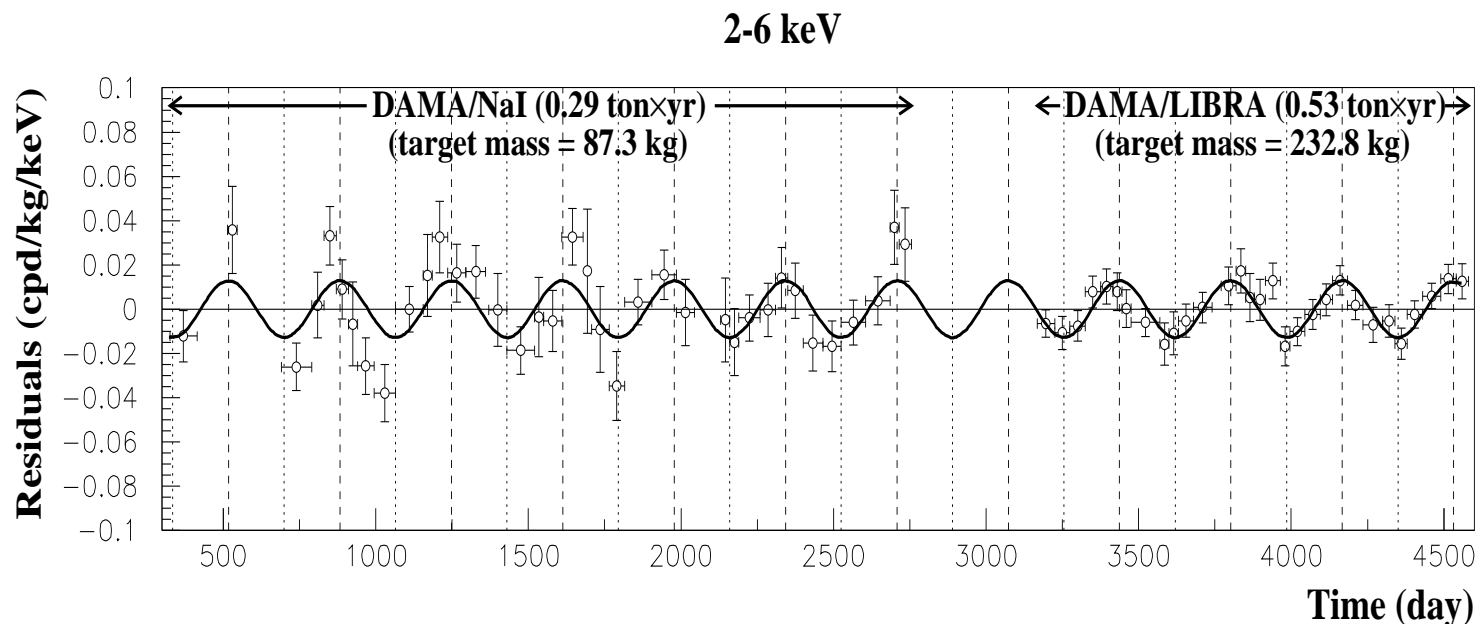
[CDMS '08]

# Annual Modulation - DAMA

- The DM flux changes due to the annual motion of the Earth relative to the galactic rotation.

[Drukier, Freese, Spergel '86]

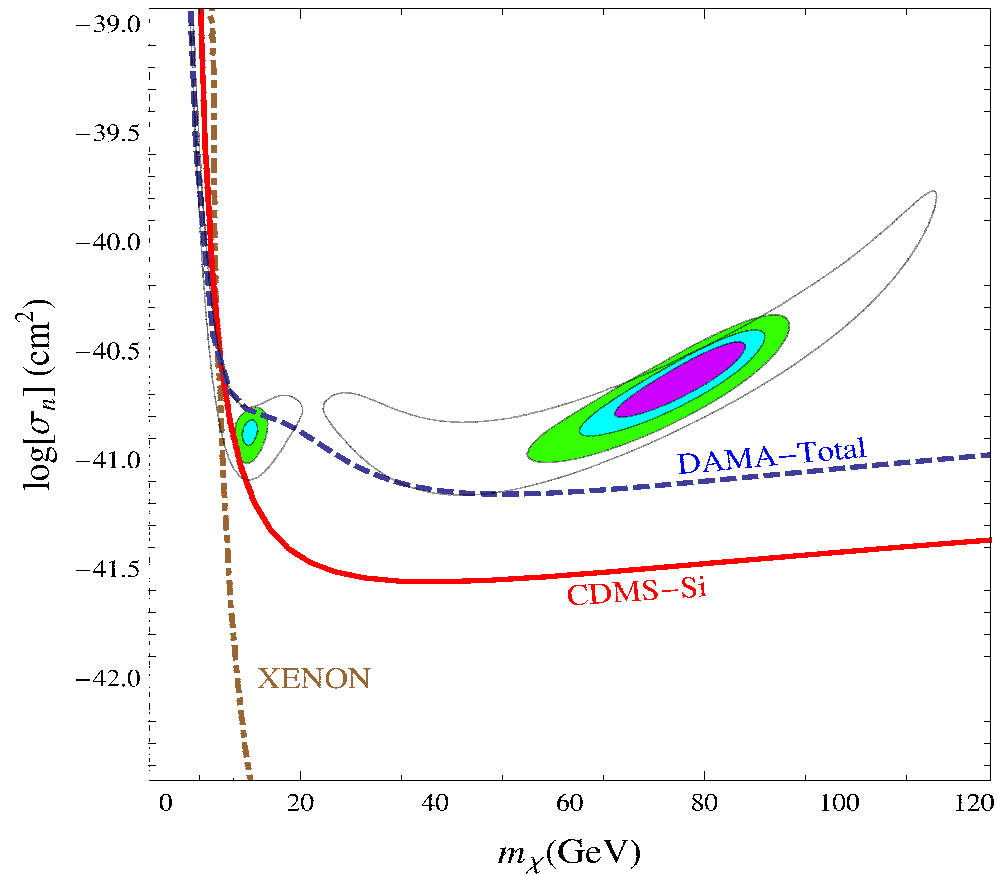
- DAMA looked for scattering off  $NaI$  with little background rejection to find an annual modulation signal ( $\gtrsim 8\sigma$ ).



# DAMA vs. Other Searches

- DAMA seems to be inconsistent with CDMS, . . . .

[Chang,Pierce,Weiner '08]



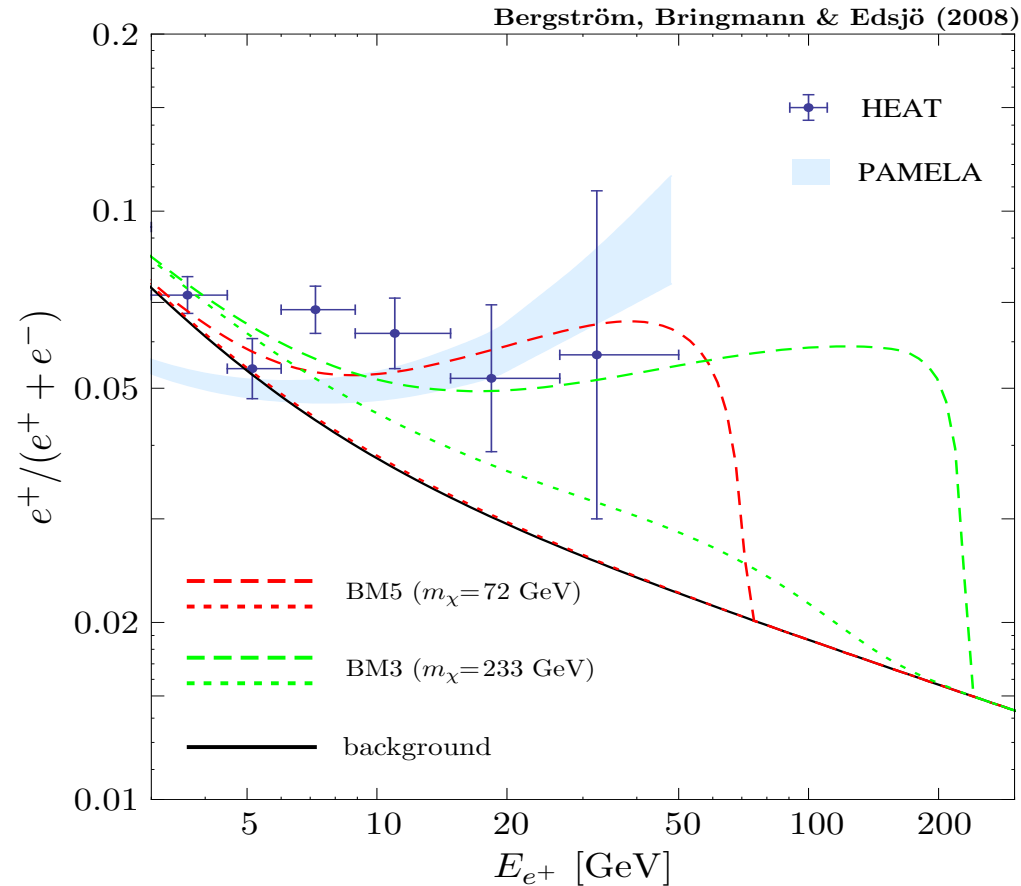
- Non-standard (inelastic) DM can still fit. [Tucker-Smith+Weiner '08]

# Indirect Detection

- Dark matter annihilation in our galaxy can give rise to particle fluxes, especially from the galactic center..
- Possible Hints:
  - PAMELA, ATIC, PPB-BETS see an excess  $e^+$  and  $e^-$ .
  - WMAP Haze: too much synchrotron radiation from the galactic center. [Finkbeiner '04]
  - INTEGRAL 511 keV line from around the galactic center.
- Signals from DM annihilating into leptons?

# Excess Cosmic Ray Positrons and Electrons

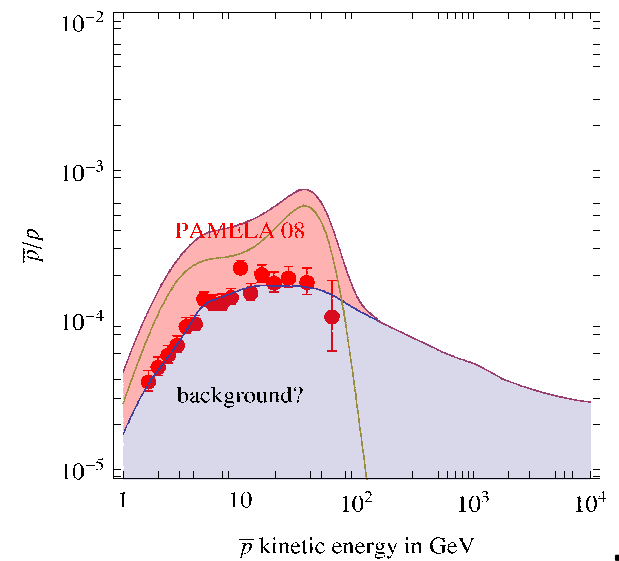
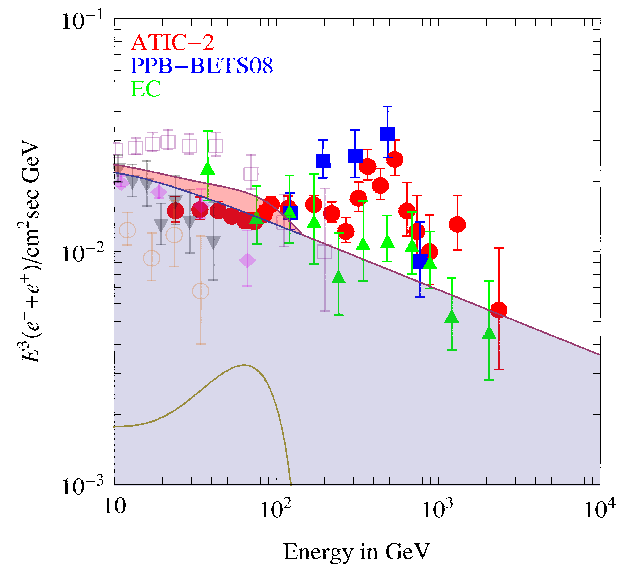
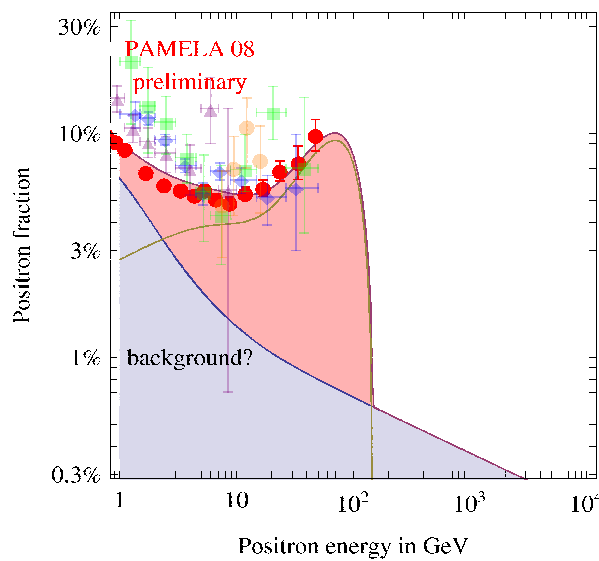
- PAMELA observes a cosmic ray positron excess.



- ATIC, PPB-BETS also see an excess  $(e^+ + e^-)$  flux.

- Energies are consistent with DM of mass  $\gtrsim 500$  GeV.
- Plots from [Cirelli, Kadastik, Raidal, Strumia '08]:  
(Preliminary Data!!!)

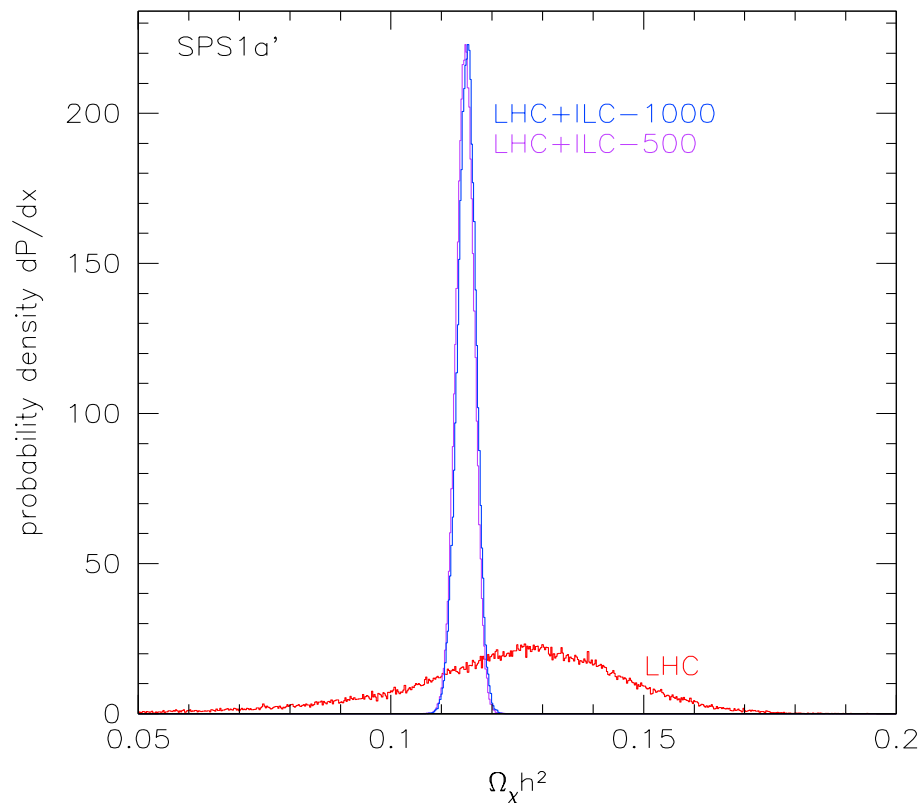
DM with  $M = 150$  GeV that annihilates into  $W^+W^-$



- Injected electrons will circulate in the galactic magnetic field and emit synchrotron radiation.  
→ **WMAP haze** [Finkbeiner '04]
- Soft  $e^+$  injected near the galactic center will annihilate.  
→ **INTEGRAL 511 keV line** [Hooper *et al.* '04]
- GLAST/Fermi should test these hints by looking for hard photons from inverse Compton scattering off the CMB.

# DM at the LHC

- DM from new physics stabilizing the electroweak scale?
- It may be possible to deduce the DM mass and couplings from data.



[Baltz *et al.* '06]

# Reconstructing the Relic Density

- If these various probes allow us to deduce the DM mass and couplings we can reconstruct  $(\Omega h^2)_{obs}$ .  
→ DM relic density assuming a standard thermal history.
- If  $(\Omega h^2)_{obs} \neq \Omega_{DM} h^2$  we learn about cosmology.
- $(\Omega h^2)_{obs} > \Omega_{DM} h^2$   
→ dilution by entropy production
- $(\Omega h^2)_{obs} < \Omega_{DM} h^2$   
→ non-thermal production, another source of DM, ...
- Another possibility for either direction is **Changing DM**.

# Changing Dark Matter

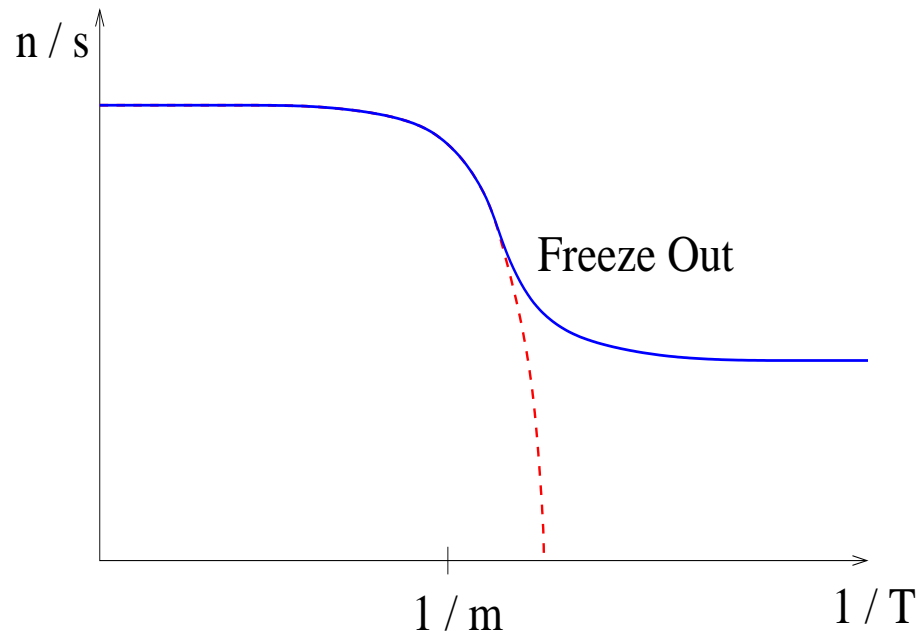
# Changing DM

- All these probes measure DM today.
- What if the DM properties changed since the time its density was formed?
- Idea: a late phase transition that modifies the DM.
- Does this have any generic features?

Is it even possible?

# WIMP Thermal Freeze-Out

- $n_{DM} \propto e^{-m/T}$  for  $T < m$  in equilibrium.
- $\Gamma_{ann} = \langle \sigma v \rangle_{ann} n_{DM}$  falls below  $H^{-1}$  at  $T = T_{fo}$ .  
 $\Rightarrow n_{DM}/s$  remains constant for  $T < T_{fo}$



# DM-Changing Phase Transition after Freeze-Out?

- Suppose  $P$  develops a VEV at  $T_{PT} < T_{fo}$ ,  
 $v_P = \langle P \rangle$  changes the properties of the DM.
- $T_{fo} \lesssim m_{DM}/20$  for WIMP DM.

We typically want  $v_P \sim m_{DM}$  for a large effect.

$$\Rightarrow T_{PT} < T_{fo} \ll v_P \sim m_{DM}$$

- Electroweak PT:  $v \sim T_{PT}$  with  $T_{PT}$  too large.
- QCD PT:  $\Lambda_{QCD} \sim T_{PT}$  with  $\Lambda_{QCD}$  too small.

# A New Phase Transition Sector (Toy Model)

- New singlet field  $P$  with potential

$$V_P = -\frac{1}{2}m_P^2 P^2 + \frac{\lambda_P}{4!}P^4$$

$$\Rightarrow v_P = \sqrt{6m_P^2/\lambda_P}$$

- $N_Q$  sets of massless fermions with  $-\mathcal{L} \supset \lambda_{PQ} P \bar{Q}Q$ .
- At finite temperature ( $P \sim 0$ ),

$$V_P \simeq -\frac{1}{2} \left( m_P^2 P^2 - \frac{N_Q}{6} \lambda_{PQ} T^2 \right) + \frac{\lambda_P}{4!} P^4$$

- The  $Q$ 's trap  $P$  at the origin independently of  $\lambda_P$ .

- At  $T = 0$

$$v_P = \sqrt{\frac{6m_P^2}{\lambda_P}}.$$

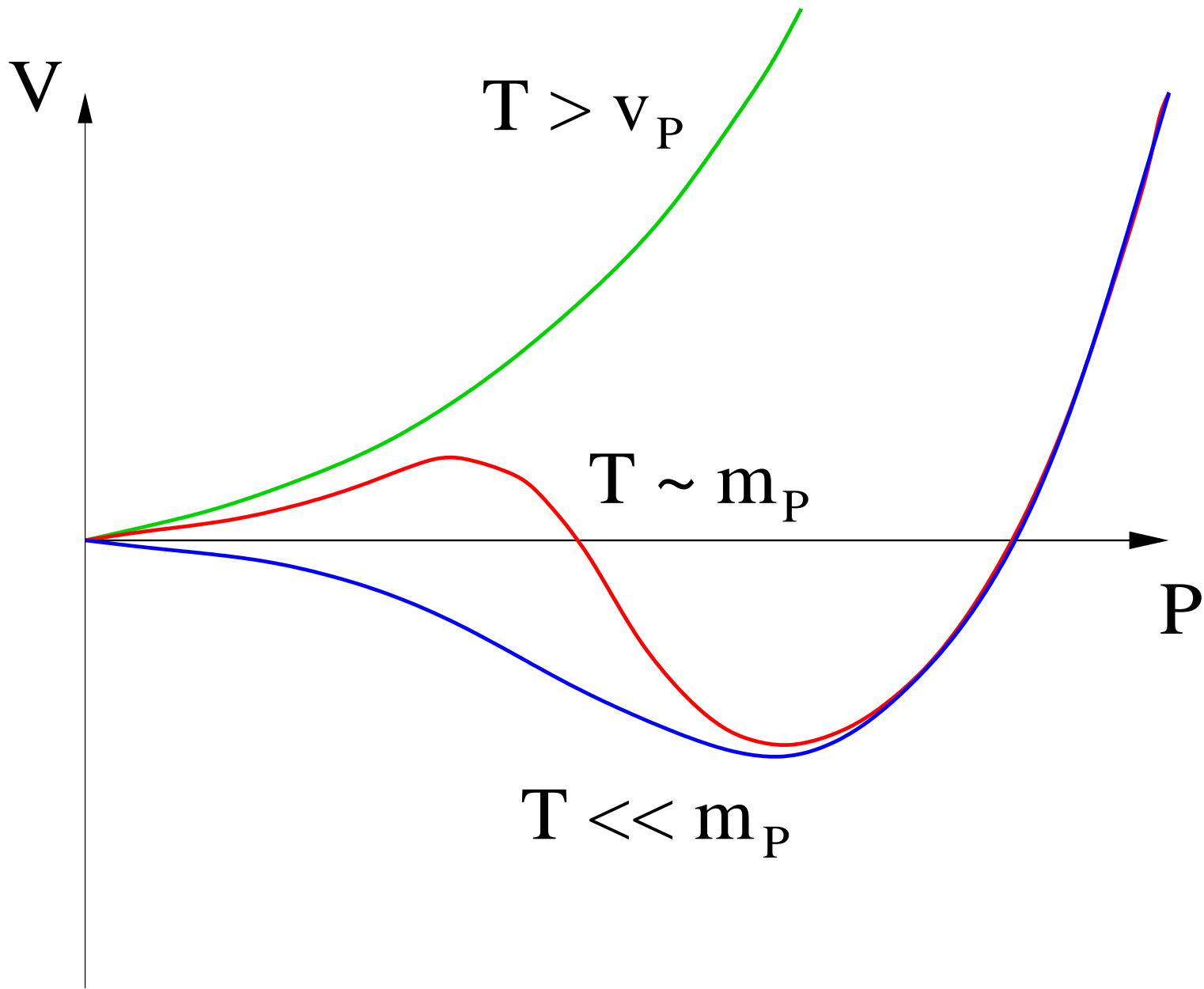
- The phase transition occurs at

$$T_{PT} = \sqrt{\frac{6m_P^2}{N_Q \lambda_{PQ}^2}}.$$

- A hierarchy  $v_P \gg T_{PT}$  arises for  $N_Q \lambda_{PQ}^2 \gg \lambda_P$ .

- Note that  $Q$  loops induce

$$\Delta\lambda_P \sim N_Q \frac{\lambda_{PQ}^4}{16\pi^2}.$$



# The Final DM Relic Density

- If this PT changes the DM mass and couplings after freeze-out, the resulting relic density is

$$\Omega_{DM} h^2 = D \left( \frac{m_{DM}^{v_P \neq 0}}{m_{DM}^{v_P = 0}} \right) \Omega_{DM}^{v_P = 0} h^2.$$

$\Omega_{DM} h^2$  is the actual value of the relic density.

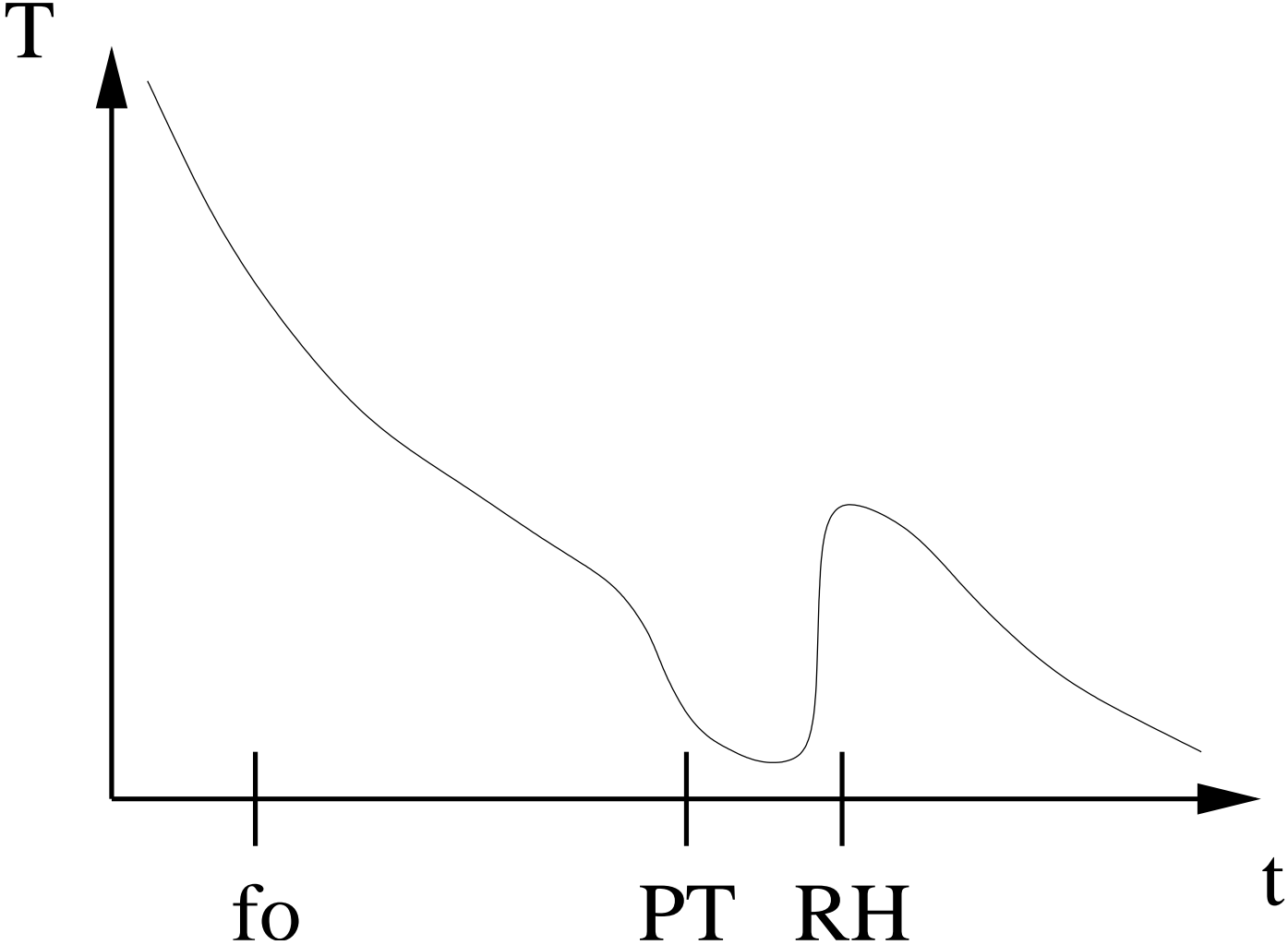
$\Omega_{DM}^{v_P = 0} h^2$  is the value obtained without a PT.

$D \leq 1$  is a dilution factor from *thermal inflation*.

$$D = \frac{g_*^{PT} T_{PT}^3}{g_*^{RH} T_{RH}^3}$$

(We focus on cases where  $D \sim 1$ .)

# Temperature Timeline



# A Sample Dark Matter Sector

- Higgsino-Singlino system with a  $\mathbb{Z}_2$  symmetry:

$SU(2)_L$  chiral doublets:  $\Psi_L, \Psi_{\bar{L}}$

Singlet chiral fermion:  $\Psi_s$

$$-\mathcal{L} \subset \lambda_s P \Psi_s \Psi_s$$

- Neutral Mass Matrix:

$$\mathcal{M} = \begin{pmatrix} 0 & \mu & -\lambda_1 v_H / \sqrt{2} \\ \mu & 0 & \lambda_2 v_H / \sqrt{2} \\ -\lambda_1 v_H / \sqrt{2} & \lambda_2 v_H / \sqrt{2} & 2(\mu_s + \lambda_s v_P) \end{pmatrix}$$

- Choose parameters such that the lightest state is:

$v_P = 0 \Rightarrow$  mostly doublet

$v_P \neq 0 \Rightarrow$  mostly singlet

# Actual and Apparent Relic Densities

- “Higgsino” DM annihilates into  $W^+W^-$  with

$$\Omega_{DM}^{v_P=0} h^2 \simeq (0.1) \left( \frac{m_{DM}^{v_P=0}}{\text{TeV}} \right)^2 .$$

- Actual relic density after the phase transition:

$$\Omega_{DM} h^2 \simeq (0.1) \left( \frac{m_{DM}^{v_P=0}}{\text{TeV}} \right)^2 D \left( \frac{m_{DM}^{v_P \neq 0}}{m_{DM}^{v_P=0}} \right) .$$

- The DM particle at  $T = 0$  is mostly singlet.
- If the DM properties at  $T = 0$  were measured, the apparent annihilation mode would be through an  $s$ -channel  $P$  to  $Q\bar{Q}$ :

$$(\Omega h^2)_{obs} \simeq \frac{0.05}{N_Q (\lambda_{PQ} \lambda_s)^2} \left( \frac{m_{DM}^{v_{P \neq 0}}}{\text{TeV}} \right)^2 .$$

- It is easy to obtain  $(\Omega h^2)_{obs} \gg \Omega_{DM} h^2$ .
- $s$ -channel  $P$  exchange also dominates the rate of annihilation in the galaxy today.

- Getting  $(\Omega h^2)_{obs} \ll \Omega_{DM} h^2$  is more difficult.
- Need  $\langle \sigma v \rangle^{v_P \neq 0} \gg \langle \sigma v \rangle^{v_P = 0}$ .  
 → modified  $v_P \neq 0$  DM tends to recouple after the PT
- Non-recoupling  $\Rightarrow n_{DM}^{PT} \langle \sigma v \rangle^{v_P \neq 0} \leq 1.66 (g_*^{RH})^{1/2} \frac{T_{RH}^2}{M_{Pl}}$ .
- Model-independent bound:

$$\frac{(\Omega h^2)_{obs}}{\Omega_{DM} h^2} \gtrsim \sqrt{\frac{g_*^{RH}}{g_*^{v_P=0}} \frac{T_{RH}}{T_{fo}^{v_P \neq 0}}}.$$

## Related Phenomenology

- A late phase transition generically requires a light scalar:

$$T_{PT} < T_{fo} \lesssim 50 \text{ GeV} \left( \frac{m_{DM}}{\text{TeV}} \right)$$

$$T_{PT} = \sqrt{m_P^2 / N_Q \lambda_{PQ}^2}$$

$m_P \sim$  mass of the scalar excitation  $\phi_P = P - v_P$ .

- “Higgs Portal” mixing:

$$-\mathcal{L} \supset \frac{\lambda_{PH}}{2} P^2 |H|^2$$

- Mixing Angle  $\theta$ :

$$\tan 2\theta = \frac{\lambda_{PH} v_P v_H}{\lambda_H^2 v_H^2 - \lambda_P^2 v_P^2}$$

## Bounds on $\phi_P$

- Heavier masses:  $m_P \gtrsim 9.5 \text{ GeV}$   
 $\Rightarrow \theta < 0.14$  from  $Z^0 h^0$  searches at LEP
- Intermediate masses:  $5.5 \text{ GeV} \lesssim m_P \lesssim 9.5 \text{ GeV}$   
 $\Rightarrow \theta \lesssim 0.3$  from  $\Upsilon \rightarrow \phi_P \gamma$  with  $\phi_P \rightarrow \tau \bar{\tau}$
- Lighter masses:  $m_P \lesssim 5.5 \text{ GeV}$   
 $\Rightarrow \theta \lesssim 10^{-4}$  from  $B^0$  decays ( $b \rightarrow s \phi_P$ )
- $h^0 Z^0 \rightarrow \phi_P \phi_P Z^0 \rightarrow 4b Z^0$  at the LHC?

[Carena, Han, Huang, Wagner '07]

## Other Possibilities

- Our phase transition sector appears to be fairly generic.  
Use the EW PT by modifying the Higgs sector?
- Other possibilities for the DM sector:
  - Move  $m_{DM}$  on or off of an  $s$ -channel resonance.
  - Change the mass of another particle that coannihilates with the DM.
  - ...

## Summary and Hints

- Lots of gravitational evidence for DM.

New data may be non-gravitational hints for DM:

DAMA, PAMELA, INTEGRAL, WMAP haze.

- Current and upcoming experiments stand a good chance of measuring the properties of the DM.

- DM properties can change after freeze-out.

- DM for PAMELA:  $\langle \sigma_{ann} v \rangle \simeq 10^{-24} \text{cm}^3/\text{s}$

with  $v \sim 10^{-3}$ .

DM for thermal freeze-out:  $\langle \sigma_{ann} v \rangle \simeq 3 \times 10^{-26} \text{cm}^3/\text{s}$

with  $v \simeq 0.05$ .