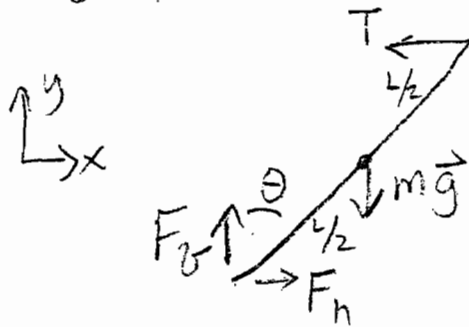


Classical Physics 171.105
 Second Mid-term Solutions 11/17/04

1. a. In equilibrium the total force and torque on the rod must vanish. Forces are drawn below.



The gravitational force is applied at the center of mass which is also the center of the stick

$$\sum \vec{F}_{\text{ext}} = 0 \Rightarrow \begin{aligned} F_h - T &= 0 && \text{(x-comp)} \\ F_v - mg &= 0 && \text{(y-comp)} \end{aligned}$$

$$F_h = T, \quad F_v = mg$$

$$\sum \tau_{\text{ext}, z} = 0 \Rightarrow T \cdot \underbrace{L \cos \theta}_{L_{\text{dist}}} - mg \underbrace{\frac{L}{2} \sin \theta}_{L_{\text{dist}}} = 0$$

use hinge as axis
but any will do

$$T = \frac{mg}{2} \tan \theta$$

b. If string cut

$$I_{zz} \alpha_z = \sum \tau_{\text{ext}, z} = -mg \frac{L}{2} \sin \theta$$

From last page case (f) $I_{zz} = \frac{1}{3} m L^2$

$$\frac{1}{3} m L^2 \alpha_z = -mg \frac{L}{2} \sin \theta$$

$$\alpha_z = -\frac{3g \sin \theta}{2L} \quad \leftarrow 0 \text{ means clockwise}$$

c. Here use conservation of energy

$$\Delta U = mg \Delta h = -mg \cdot \frac{L}{2} \cos \theta \quad \leftarrow -\frac{L}{2} \cos \theta = \text{height change of c.m.}$$

$$\Delta K = -\Delta U = mg \cdot \frac{L}{2} \cos \theta = \frac{1}{2} I \omega_f^2 \quad \text{since } \omega_i = 0$$

$$mg \cdot \frac{L}{2} \cos \theta = \frac{1}{2} \cdot \frac{1}{3} m L^2 \omega_f^2$$

$$\omega_f^2 = \frac{3g \cos \theta}{L}, \quad \omega_f = \sqrt{\frac{3g \cos \theta}{L}}$$

2. a. $\Delta U = -MgH$

$$\Delta K = \frac{1}{2}Mv_f^2 = -\Delta U \Rightarrow \frac{1}{2}Mv_f^2 = MgH$$

$$v_f = \sqrt{2gH}$$

b. The change in mechanical energy is the change in $U+K$.

$$\begin{aligned} \Delta E_{\text{mech}} &= \Delta U + \Delta K = -MgH + \frac{1}{2}M(0.7 \cdot \sqrt{2gH})^2 \\ &= -MgH [1 - 0.7^2] = -0.51MgH \end{aligned}$$

This energy went into internal energy in the baller and the hillside, included by motion of other rollers

c. If rolls without slipping must rotate with

$$|w|R = v. \text{ Now } \Delta K = \frac{1}{2}Mv_f^2 + \frac{1}{2}Iw_f^2$$

$$-\Delta U = \frac{1}{2}Mv_f^2 + \frac{1}{2} \cdot \frac{2}{5}MR^2 \cdot \left(\frac{v_f}{R}\right)^2$$

$$MgH = Mv_f^2 \left(\frac{1}{2} + \frac{1}{5}\right) = 0.7 Mv_f^2$$

$$v_f = \sqrt{\frac{10}{7}gH}$$

3a. Use conservation of angular momentum.

$$\text{Initially } \vec{L}_{\text{tot}} = 0$$

$$\text{Afterwards } 0 = \vec{L}_{\text{tot}} = \underbrace{m_2 v R}_{\substack{\vec{r} \times \vec{p} \\ \text{for rock}}} + \underbrace{I \omega}_{\text{meringgo round}}$$

$$\omega = - \frac{m_2 v R}{I}$$

$$I = \frac{1}{2} M_1 R^2$$

$$= - \frac{2m_2 \cdot v}{M_1 R}$$

sign \Rightarrow clockwise

$$|\omega| = \frac{2m_2 v}{M_1 R} \quad |\vec{L}| = m_2 v R.$$

b. $\vec{\omega}$ and \vec{L} are both into page, by right hand rule

c. linear speed's $|\omega R| = \frac{2m_2}{M_1} v$

4.a. At each r only the parts of the earth that are at smaller r contribute to the force. They act like a sphere of the same mass centered on the origin. The mass is

$$M_{er} = \underbrace{\frac{4\pi}{3} r^3}_{\text{volume}} \cdot \underbrace{\rho}_{\text{density}}$$

For $r=R$ have $\frac{4\pi}{3} R^3 \rho = M_e \Rightarrow \rho = \frac{3M_e}{4\pi R^3}$

$$M_{er} = \frac{r^3}{R^3} \cdot M_e$$

$|\vec{F}| = G \frac{M_{er} m}{r^2}$ directed toward center $|\vec{F}| = \frac{m M_e r}{G R^3}$

b. $U = - \int_0^r \vec{F} \cdot d\vec{s}$ As go away from origin $\vec{F} \cdot d\vec{s} = -F dr$ since pulled back to center

$$= + \int_0^r \frac{G M_e m r}{R^3} dr = \frac{G M_e m r^2}{2 R^3}$$

c. $\frac{1}{2} m v^2 \Big|_{\Delta R} = - \left[\frac{G M_e m}{2 R^3} R^2 - \frac{G M_e m r^2}{2 R^3} \right]$

↑
initial U

$$v^2 = \frac{G M_e}{R^3} (R^2 - r^2) \quad v = \sqrt{\frac{G M_e}{R^3} (R^2 - r^2)}$$